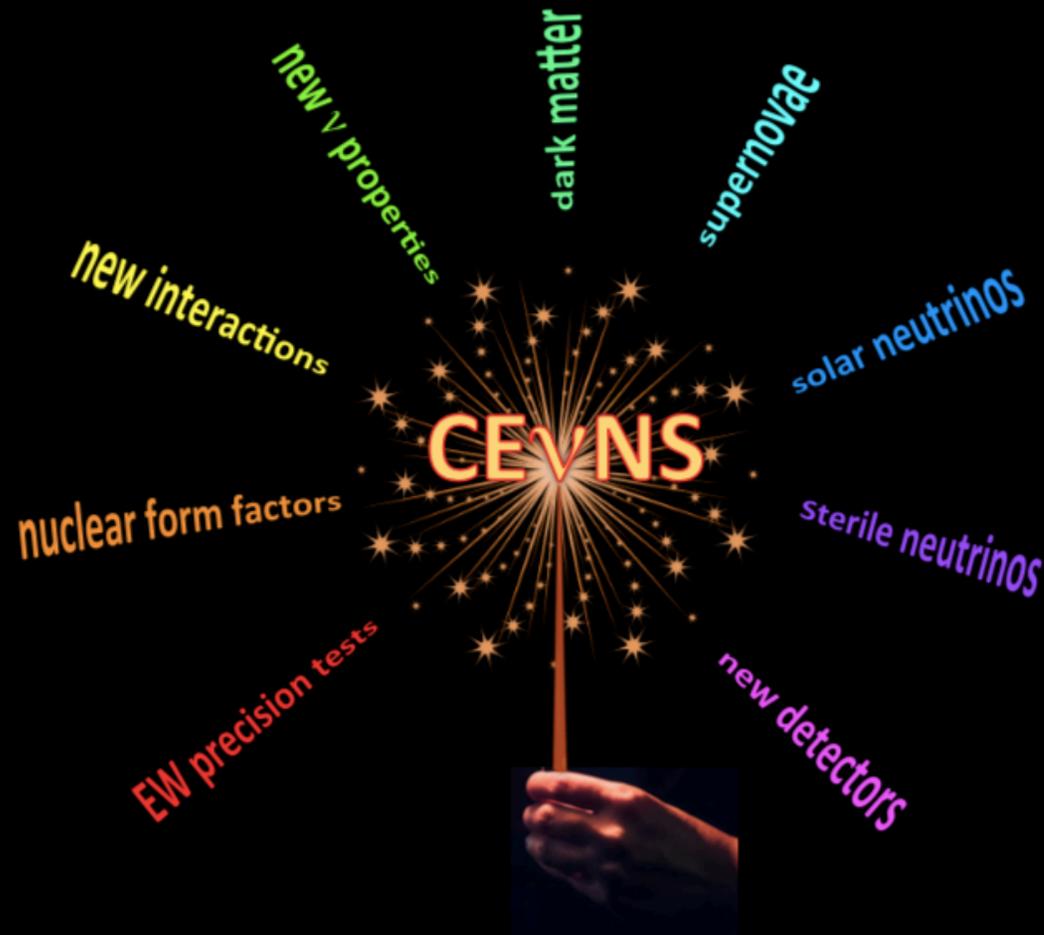
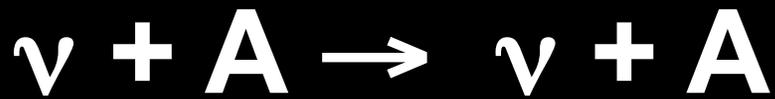


COHERENT Physics Program

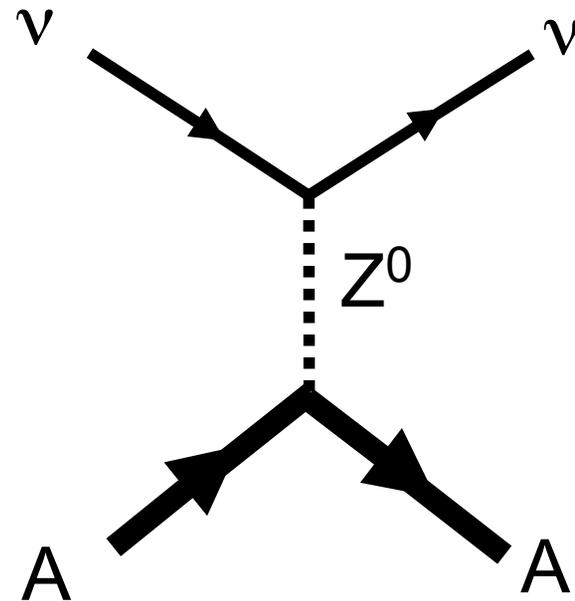
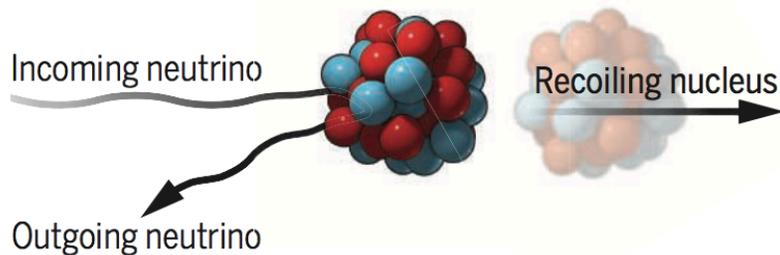


K. Scholberg
COHERENT Review
August 15, 2018

Coherent elastic neutrino-nucleus scattering (CEvNS)



A neutrino smacks a nucleus via exchange of a Z , and the nucleus recoils as a whole; **coherent** up to $E_\nu \sim 50$ MeV



Nucleon wavefunctions in the target nucleus are **in phase with each other** at low momentum transfer

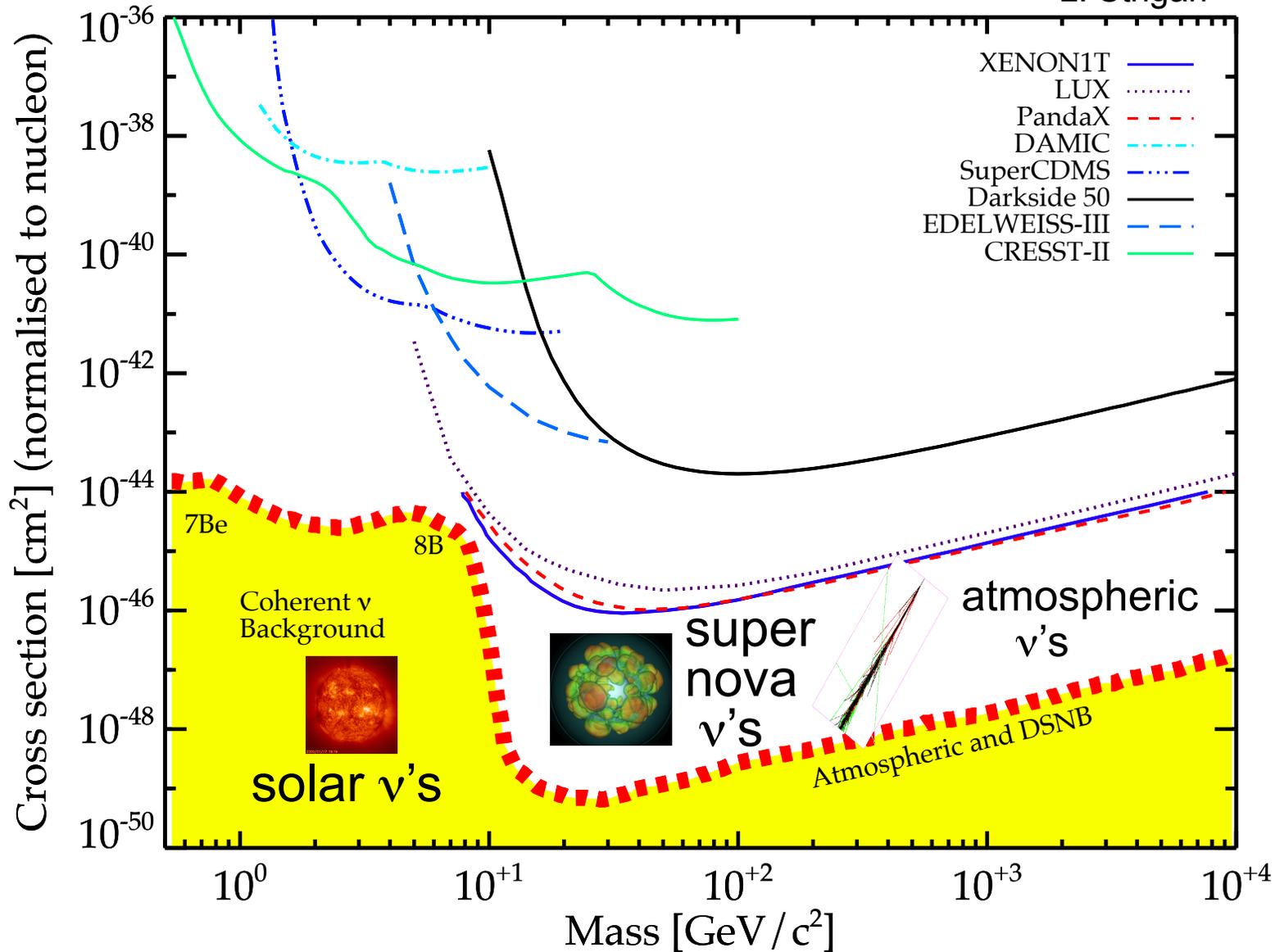
$$\text{For } QR \ll 1, \quad [\text{total xscn}] \sim A^2 * [\text{single constituent xscn}]$$

The so-called “neutrino floor” (signal!) for DM experiments

J. Monroe & P. Fisher, 2007

J. Billard, E. Figueroa-Feliciano, and L. Strigari, arXiv:1307.5458v2 (2013).

L. Strigari



Measure CEvNS to understand nature of background/astro signal
(& detector response, DM interaction)

The cross section is cleanly predicted in the Standard Model

$$\frac{d\sigma}{dT} = \frac{G_F^2 M}{\pi} F^2(Q) \left[(G_V + G_A)^2 + (G_V - G_A)^2 \left(1 - \frac{T}{E_\nu}\right)^2 - (G_V^2 - G_A^2) \frac{MT}{E_\nu^2} \right]$$

E_ν : neutrino energy

T: nuclear recoil energy

M: nuclear mass

$Q = \sqrt{2 M T}$: momentum transfer

G_V, G_A : SM weak parameters

vector

$$G_V = g_V^p Z + g_V^n N,$$

← dominates

axial

$$G_A = g_A^p (Z_+ - Z_-) + g_A^n (N_+ - N_-)$$

← small for
most
nuclei,
zero for
spin-zero

$$g_V^p = 0.0298$$

$$g_V^n = -0.5117$$

$$g_A^p = 0.4955$$

$$g_A^n = -0.5121.$$

The cross section is cleanly predicted in the Standard Model

$$\frac{d\sigma}{dT} = \frac{G_F^2 M}{\pi} F^2(Q) \left[(G_V + G_A)^2 + (G_V - G_A)^2 \left(1 - \frac{T}{E_\nu}\right)^2 - (G_V^2 - G_A^2) \frac{MT}{E_\nu^2} \right]$$

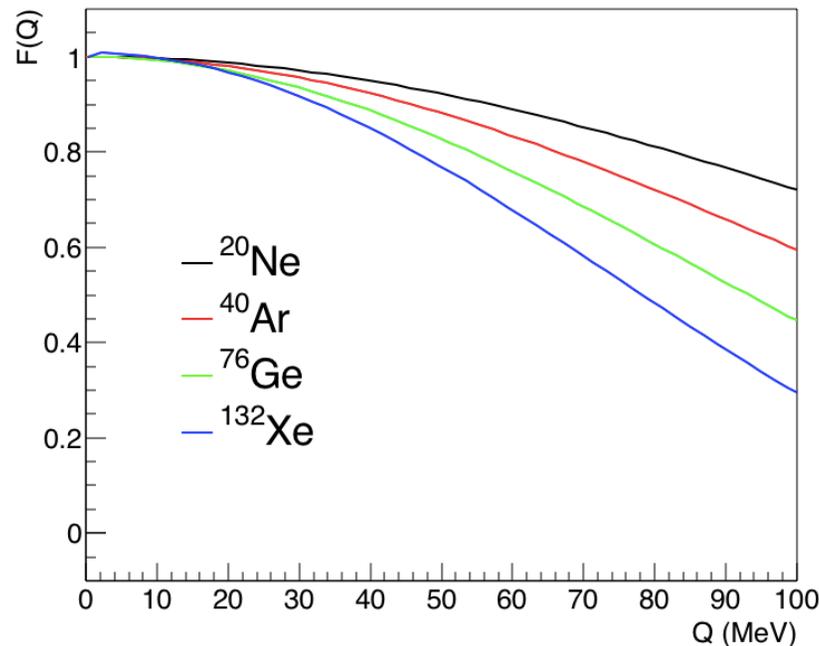
E_ν : neutrino energy

T : nuclear recoil energy

M : nuclear mass

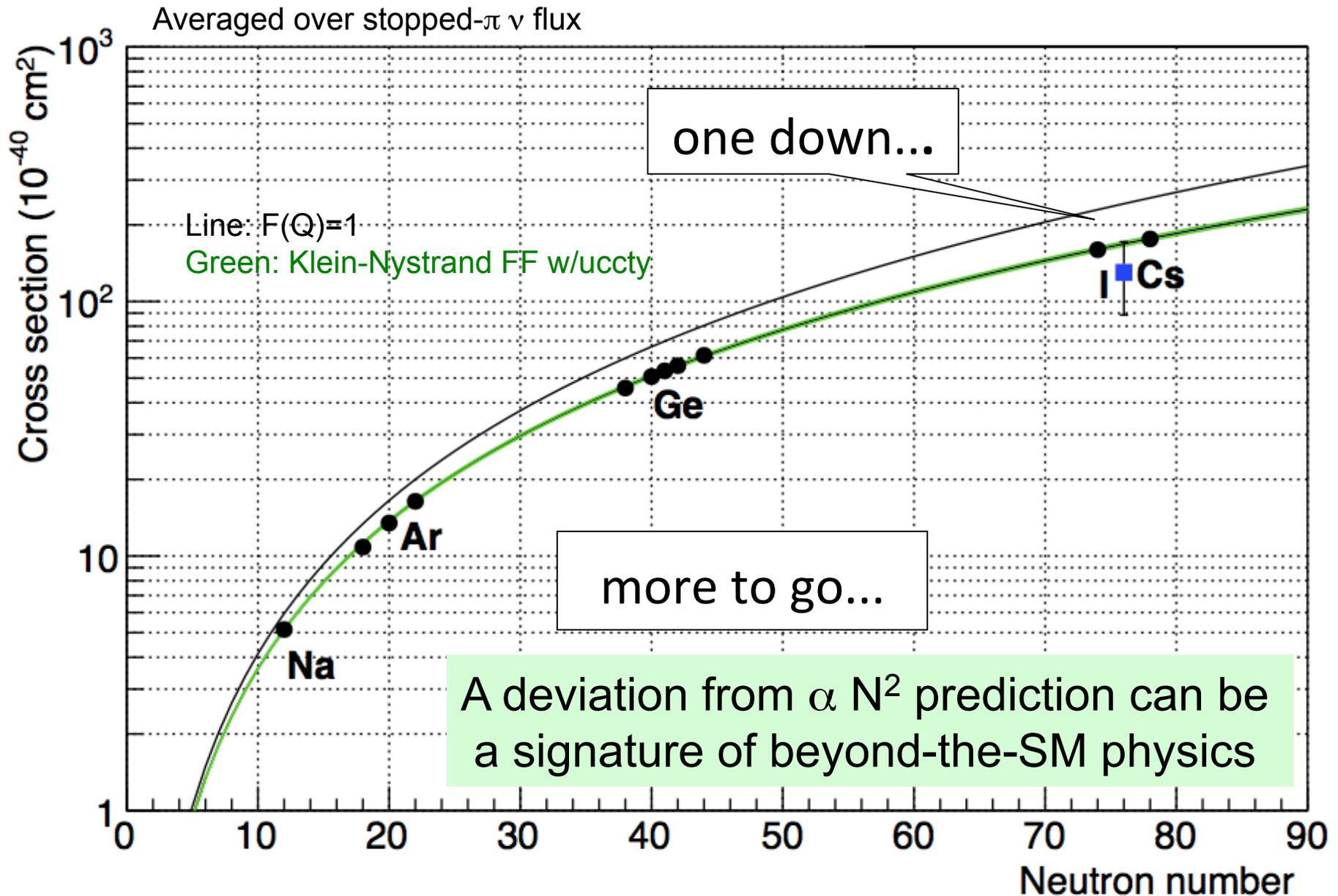
$Q = \sqrt{2 M T}$: momentum transfer

$F(Q)$: nuclear **form factor**, $<\sim 5\%$ uncertainty on event rate

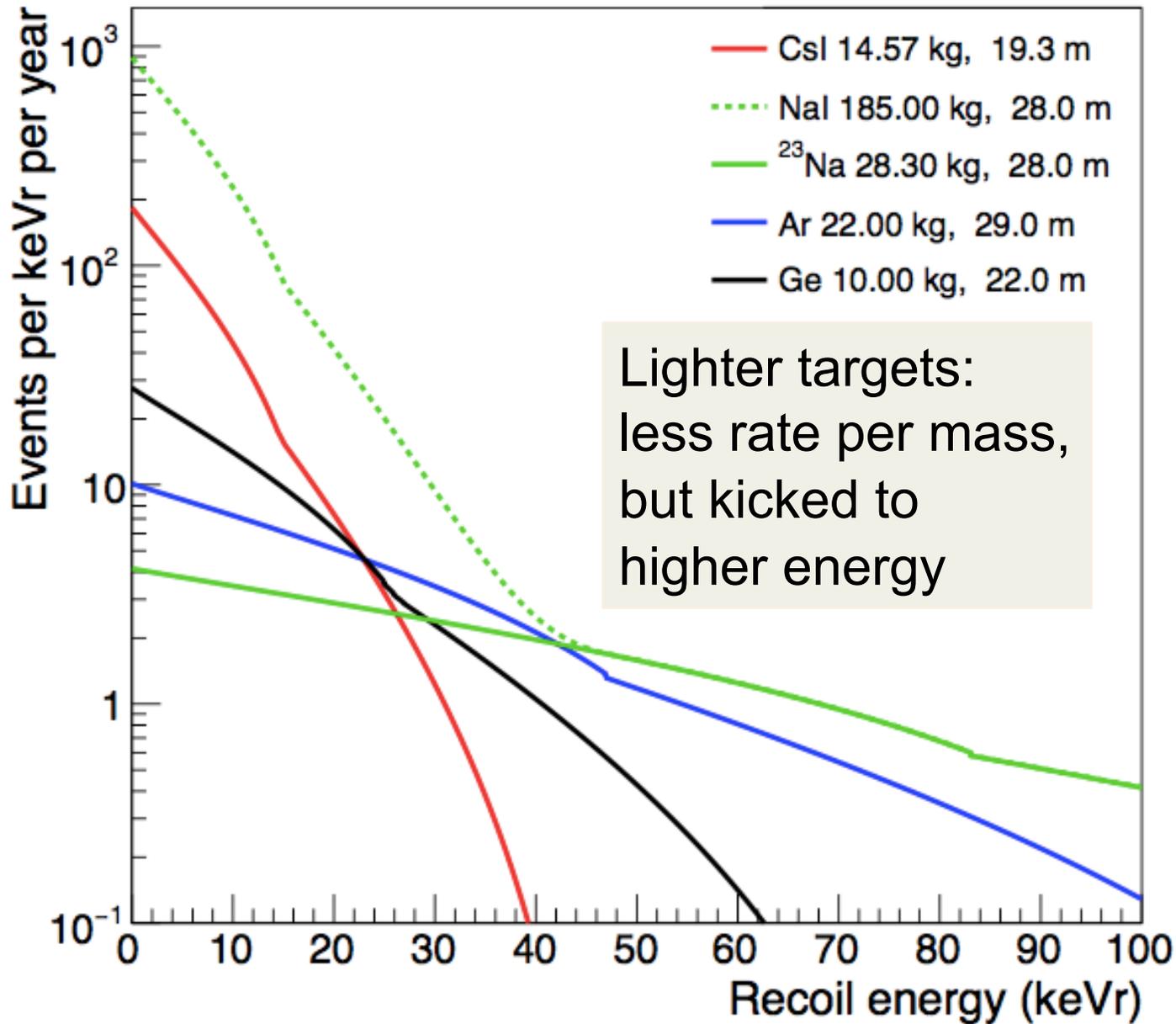


form factor
suppresses
cross section
at large Q

Need to measure N^2 dependence of the CEvNS xscn



COHERENT's targets: expected recoil energy distributions



COHERENT Physics Topics

Topic	Experimental Signature
Non-standard-interactions/new mediators	Deviation from N^2 , deviation from SM recoil spectrum shape
Weak mixing angle measurement	Event rate scaling
Neutrino magnetic moment	Excess at low recoil energy
Accelerator-produced DM	Event rate scaling
Sterile oscillations	Event rate and spectrum at multiple baselines
Nuclear form factors	Recoil spectrum shape
Inelastic CC/NC xscn for supernova	High-energy (MeV) electrons, deex γ 's
Inelastic CC/NC xscn for weak coupling parameters	High-energy (MeV) electrons, deex γ 's

COHERENT Physics Topics

Topic	Experimental Signature	Detector Requirements
Non-standard-interactions/new mediators	Deviation from N^2 , deviation from SM recoil spectrum shape	Multiple targets, energy resolution, QF
Weak mixing angle measurement	Event rate scaling	Multiple targets, QF
Neutrino magnetic moment	Excess at low recoil energy	Low energy threshold and resolution, QF
Accelerator-produced DM	Event rate scaling	Energy resolution, QF
Sterile oscillations	Event rate and spectrum at multiple baselines	Similar/movable detectors at multiple baselines
Nuclear form factors	Recoil spectrum shape	Energy resolution, QF
Inelastic CC/NC xscn for supernova	High-energy (MeV) electrons, deex γ 's	Large target mass
Inelastic CC/NC xscn for weak coupling parameters	High-energy (MeV) electrons, deex γ 's	Large target mass

Physics topics by detector subsystem

	Csl	Ar	NaI	Ge	Nubes	D ₂ O
Non-standard-interactions/ new mediators	✓	✓	✓	✓		
Weak mixing angle measurement	✓	✓	✓	✓		
Neutrino magnetic moment				✓		
Accelerator-produced DM	✓	✓	✓	✓		
Sterile oscillations	✓	✓	✓	✓		
Nuclear form factors	✓	✓	✓	✓		
Inelastic CC/NC xscn for supernova		✓			✓	✓
Inelastic CC/NC xscn for weak coupling parameters		✓	✓		✓	

Grey texture: **combination important** (N^2 , systematics cancellation)

D₂O important for all

Quenching factor measurements critical for all

Searching for BSM Physics with CEvNS

A first example: simple counting to constrain
non-standard interactions (NSI) of
neutrinos with quarks

Davidson et al., JHEP 0303:011 (2004)
Barranco et al., JHEP 0512:021 (2005)

“Model-independent” parameterization

$$\mathcal{L}_{\nu H}^{NSI} = -\frac{G_F}{\sqrt{2}} \sum_{\substack{q=u,d \\ \alpha,\beta=e,\mu,\tau}} [\bar{\nu}_\alpha \gamma^\mu (1 - \gamma^5) \nu_\beta] \times (\varepsilon_{\alpha\beta}^{qL} [\bar{q} \gamma_\mu (1 - \gamma^5) q] + \varepsilon_{\alpha\beta}^{qR} [\bar{q} \gamma_\mu (1 + \gamma^5) q])$$

ε 's parameterize new interactions

“Non-Universal”: ε_{ee} , $\varepsilon_{\mu\mu}$, $\varepsilon_{\tau\tau}$

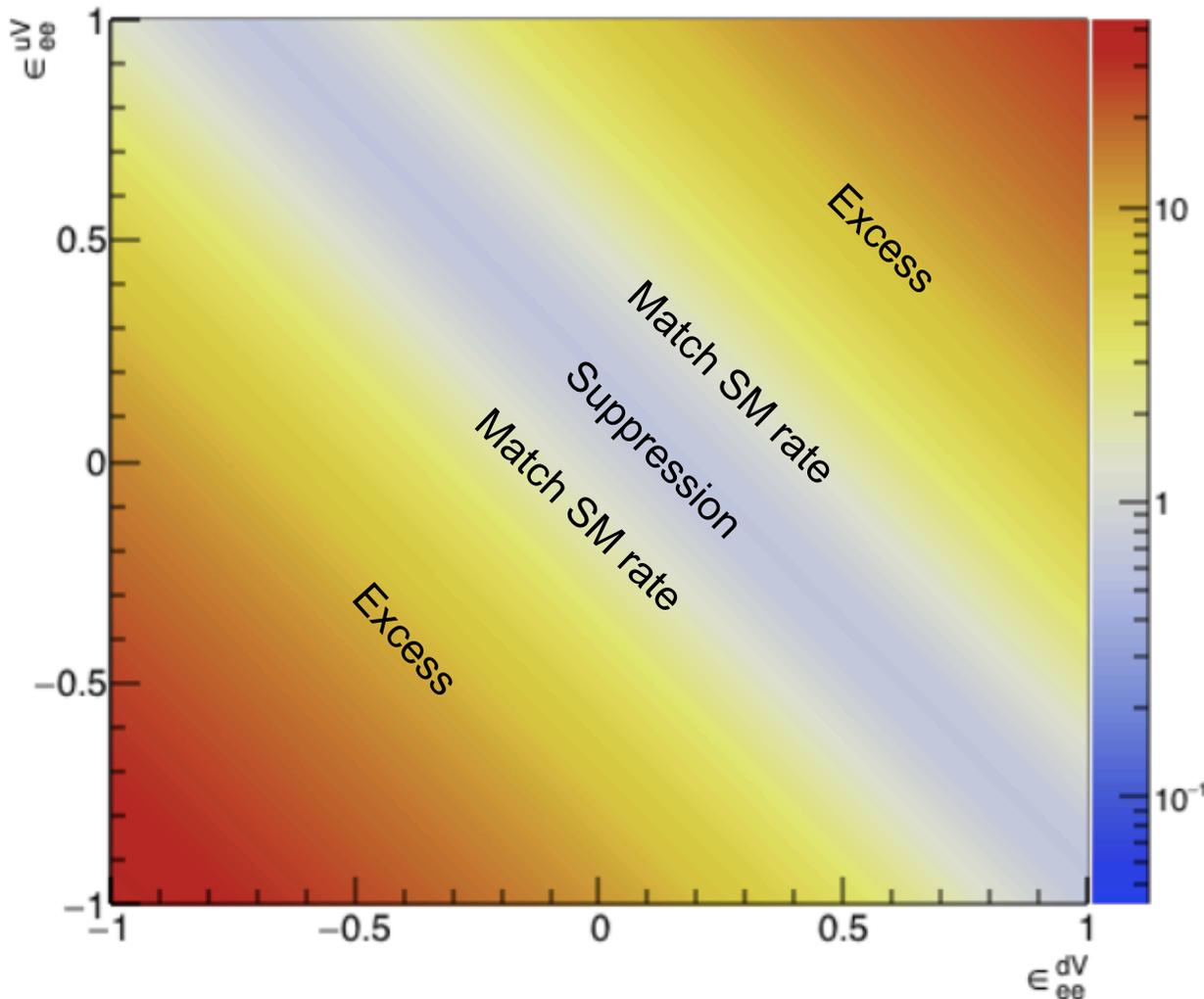
Flavor-changing: $\varepsilon_{\alpha\beta}$, where $\alpha \neq \beta$

\Rightarrow some are quite poorly constrained (\sim unity allowed)

Ratio of rate with NSI to SM rate (all flavors in stopped-pion beam)

ϵ_{ee}^{uV} vs ϵ_{ee}^{dV} parameters (assume others zero)

Csl



Note that for

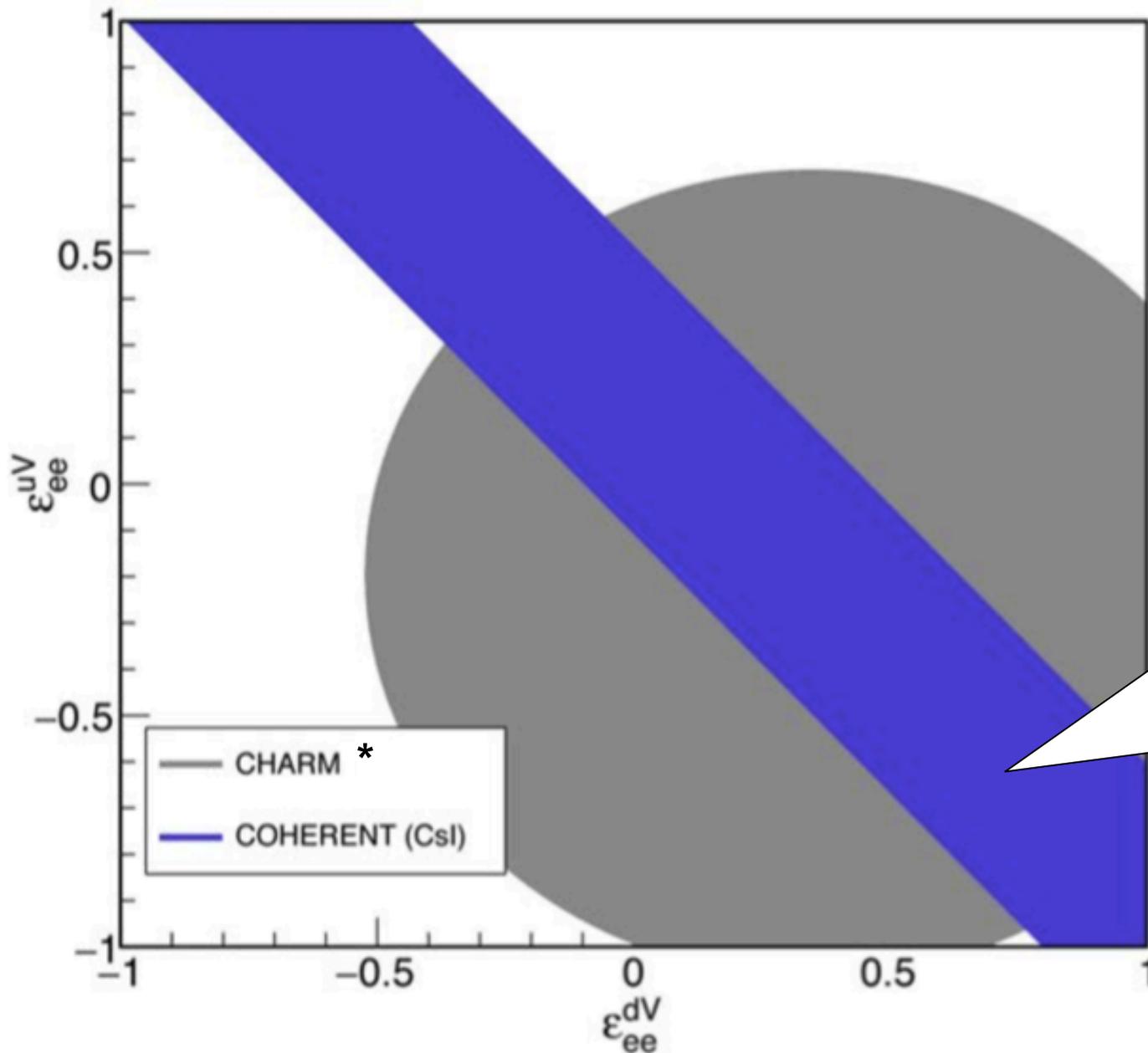
$$Z(g_V^p + 2\epsilon_{ee}^{uV} + \epsilon_{ee}^{dV}) + N(g_V^n + \epsilon_{ee}^{uV} + 2\epsilon_{ee}^{dV}) = \pm(Zg_V^p + Ng_V^n),$$

the rate is the same as for the SM, so parameters will be allowed

Get slightly different slope for different targets

Example: Neutrino non-standard interactions

Constraints for current CsI data set:



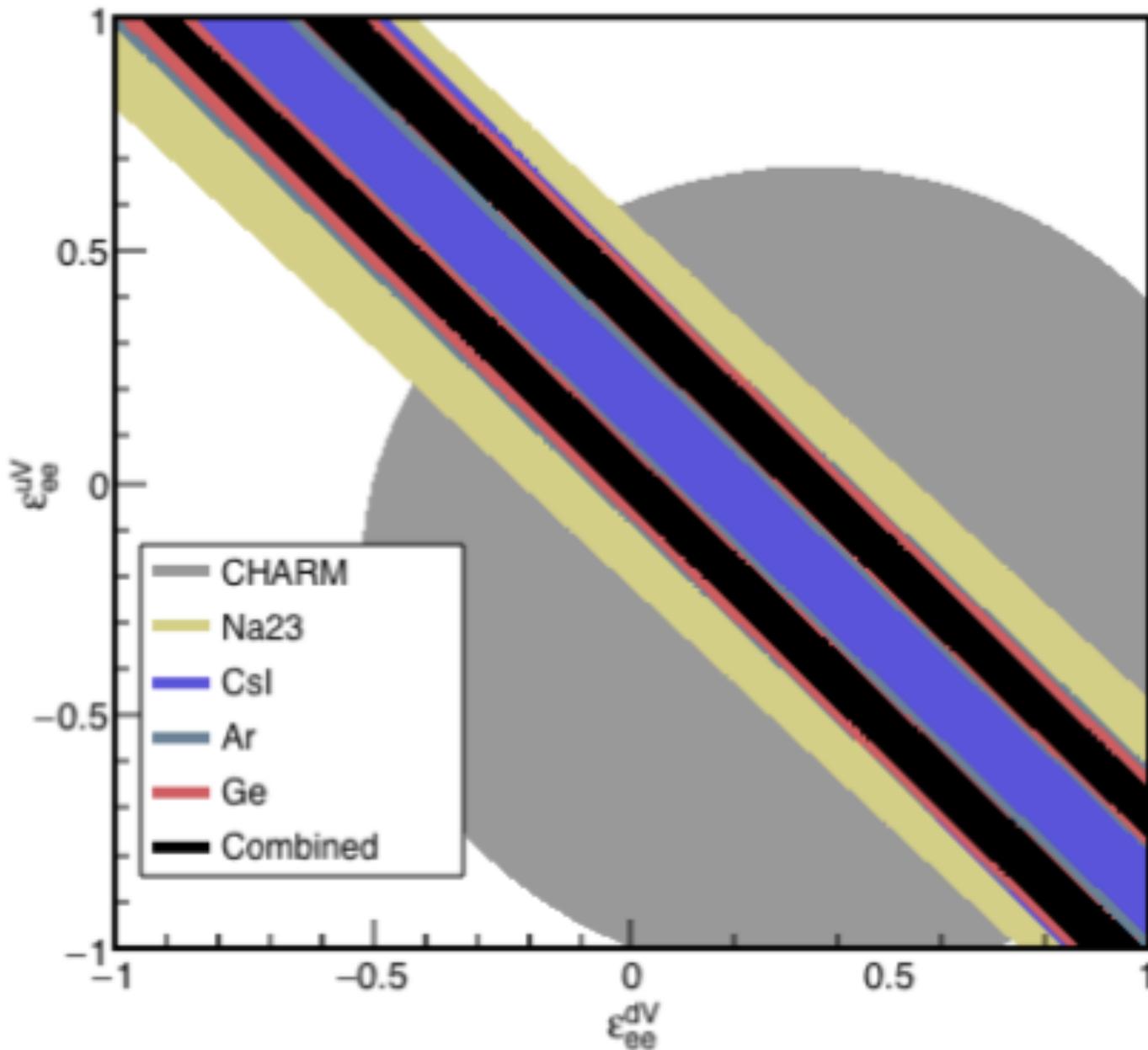
- Assume all other ϵ 's zero

Parameters describing beyond-the-SM interactions outside this region disfavored at 90%

See also Coloma et al., arXiv:1708.02899

*CHARM constraints apply only to heavy mediators

Projected future sensitivities for NSI

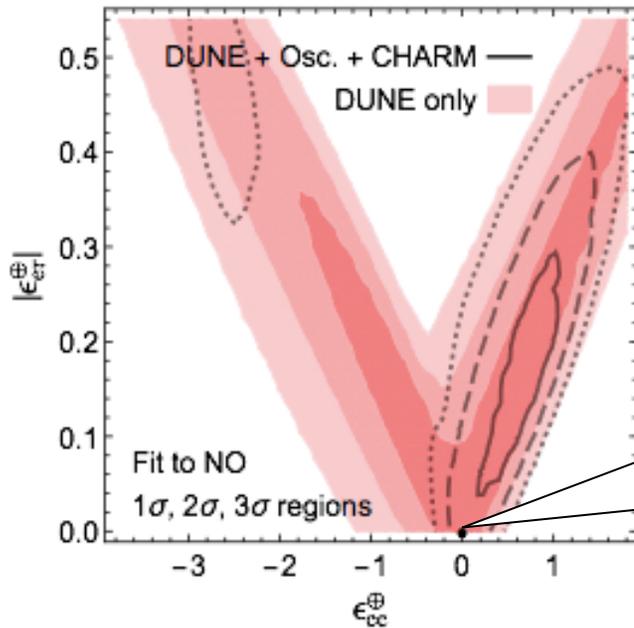


Combination
of targets
improves
sensitivity

Generalized mass ordering degeneracy in neutrino oscillation experiments

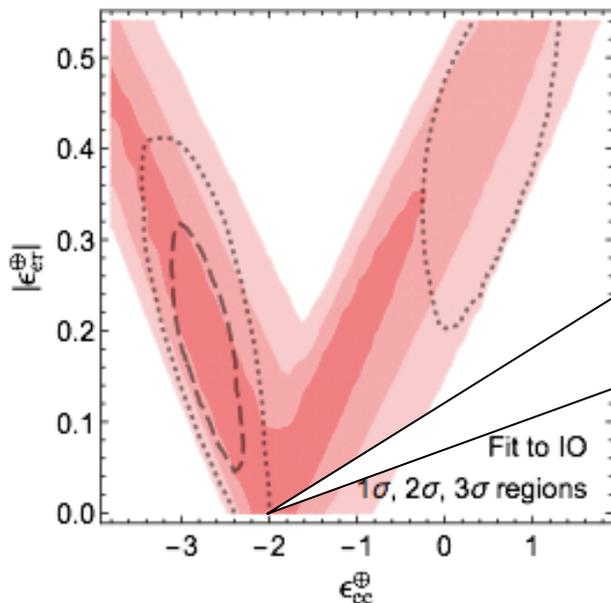
Pilar Coloma¹ and Thomas Schwetz²

Phys.Rev. D94 (2016) no.5, 055005,
 Erratum: Phys.Rev. D95 (2017) no.7, 079903
 P. Coloma et al., JHEP 1704 (2017) 116



Normal
ordering
w/no
NSI...

If you allow for NSI,
 an ambiguity
 exists in determining
 mass ordering
 w/ LBL experiments:
“LMA-Dark”



...looks
just like
inverted
ordering
w/NSI

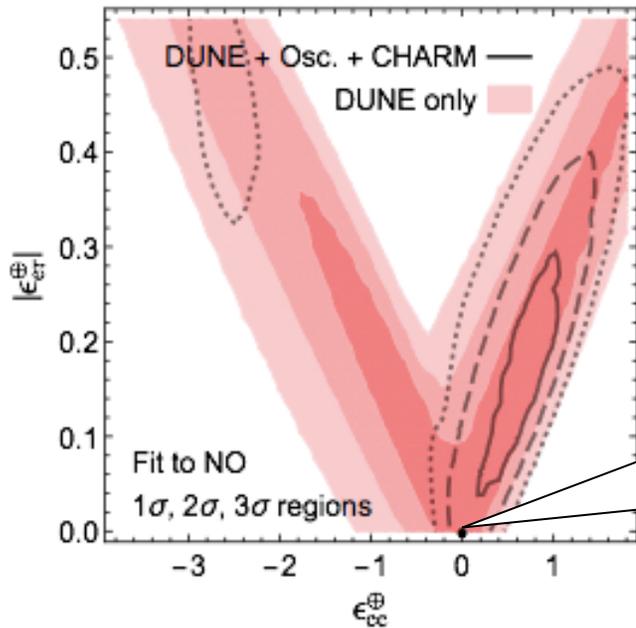
Same answer for

$$\begin{aligned} \Delta m_{31}^2 &\rightarrow -\Delta m_{31}^2 + \Delta m_{21}^2 = -\Delta m_{32}^2, \\ \sin \theta_{12} &\rightarrow \cos \theta_{12}, \\ \delta &\rightarrow \pi - \delta, \\ (\epsilon_{ee} - \epsilon_{\mu\mu}) &\rightarrow -(\epsilon_{ee} - \epsilon_{\mu\mu}) - 2, \\ (\epsilon_{\tau\tau} - \epsilon_{\mu\mu}) &\rightarrow -(\epsilon_{\tau\tau} - \epsilon_{\mu\mu}), \\ \epsilon_{\alpha\beta} &\rightarrow -\epsilon_{\alpha\beta}^* \quad (\alpha \neq \beta) \end{aligned}$$

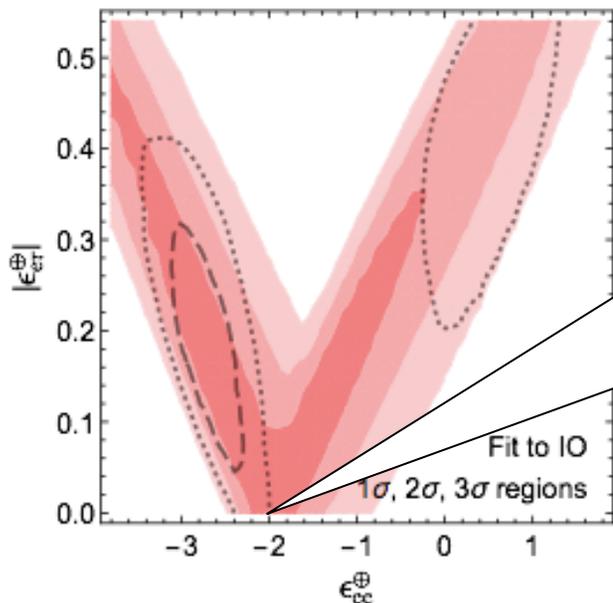
Generalized mass ordering degeneracy in neutrino oscillation experiments

Pilar Coloma¹ and Thomas Schwetz²

Phys.Rev. D94 (2016) no.5, 055005,
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P. Coloma et al., JHEP 1704 (2017) 116



Normal ordering w/no NSI...



...looks just like inverted ordering w/NSI

CEvNS measurements can place significant constraints to resolve the LMA-D ambiguity if SM rate is measured

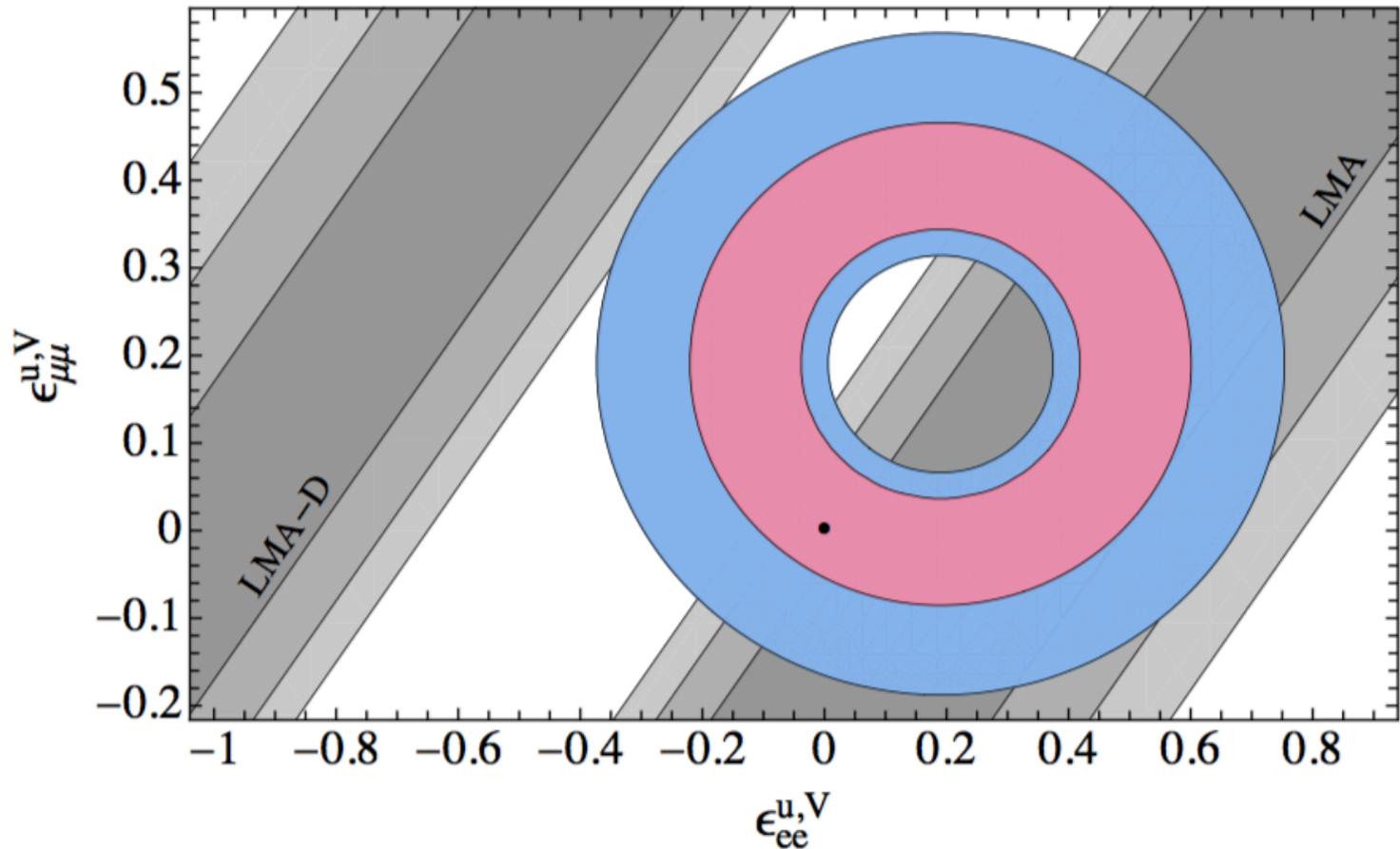
OR, could confirm an NSI signature observed by DUNE

A COHERENT enlightenment of the neutrino Dark Side

Pilar Coloma,^{1,*} M. C. Gonzalez-Garcia,^{2,3,4,†} Michele Maltoni,^{5,‡} and Thomas Schwetz^{6,§}

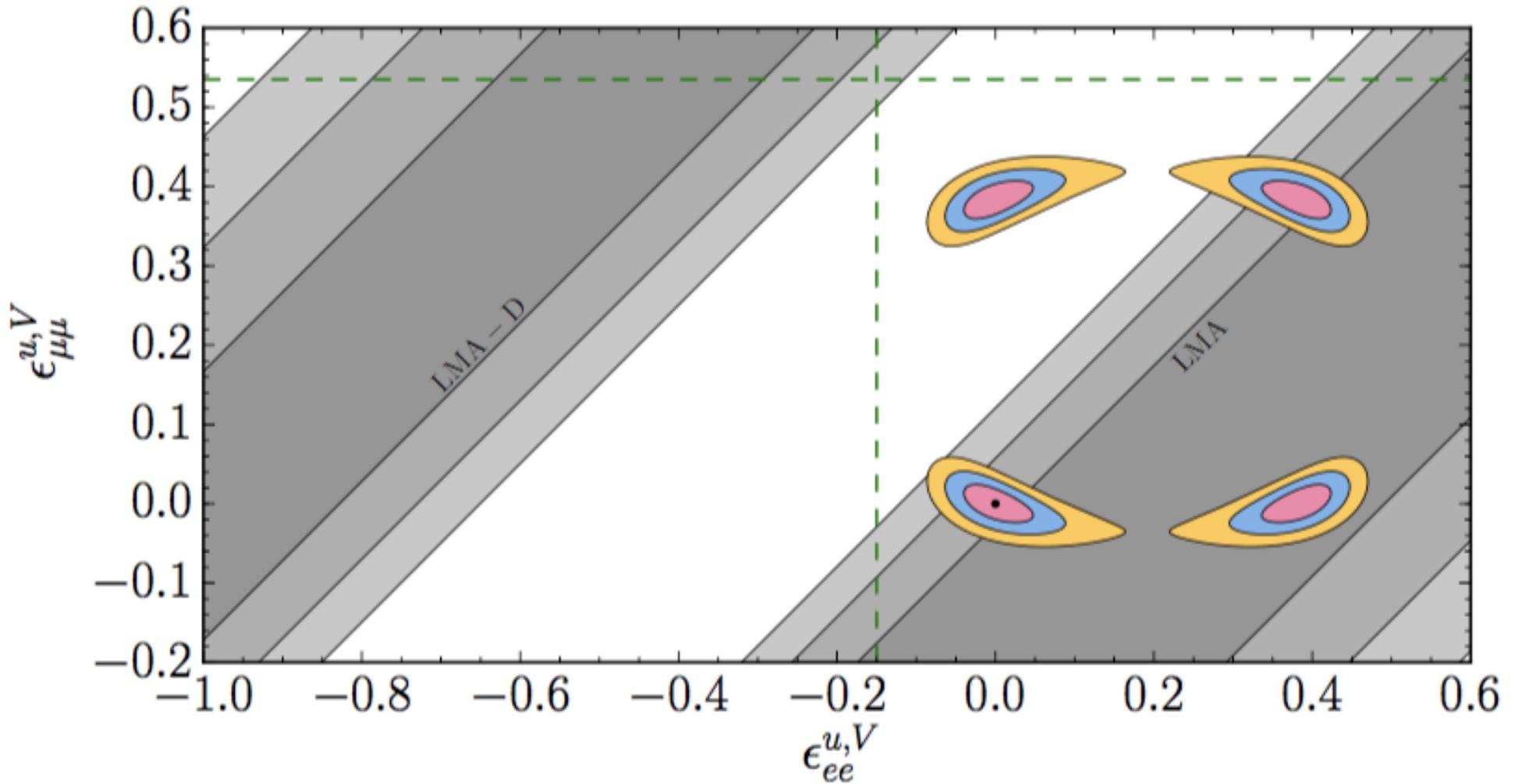
Phys.Rev. D96 (2017) no.11, 115007

1 σ , 2 σ
allowed
regions
projected in
(ϵ_{ee}^{uV} , $\epsilon_{\mu\mu}^{uV}$)
plane



First COHERENT results are already disfavoring LMA-D

Future COHERENT results will fully exclude LMA-D



Another phenomenological analysis, making use of spectral fit:

COHERENT constraints on nonstandard neutrino interactions

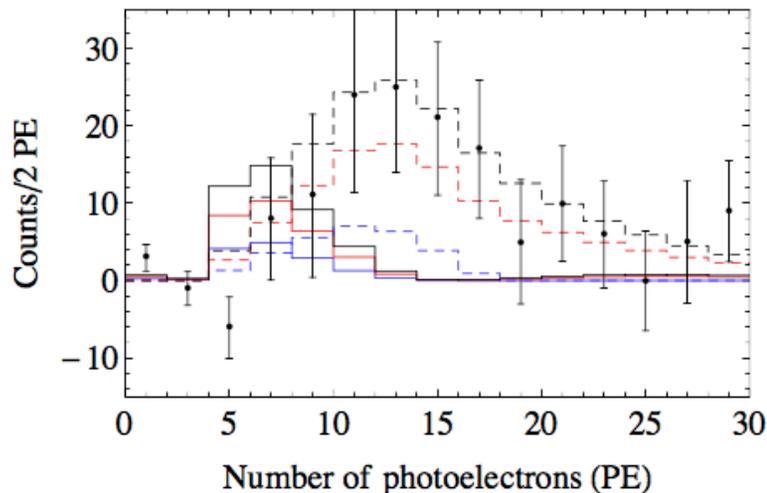
Jiajun Liao and Danny Marfatia
arXiv:1708.04255

SM weak charge

Effective weak charge in presence of light vector mediator Z'

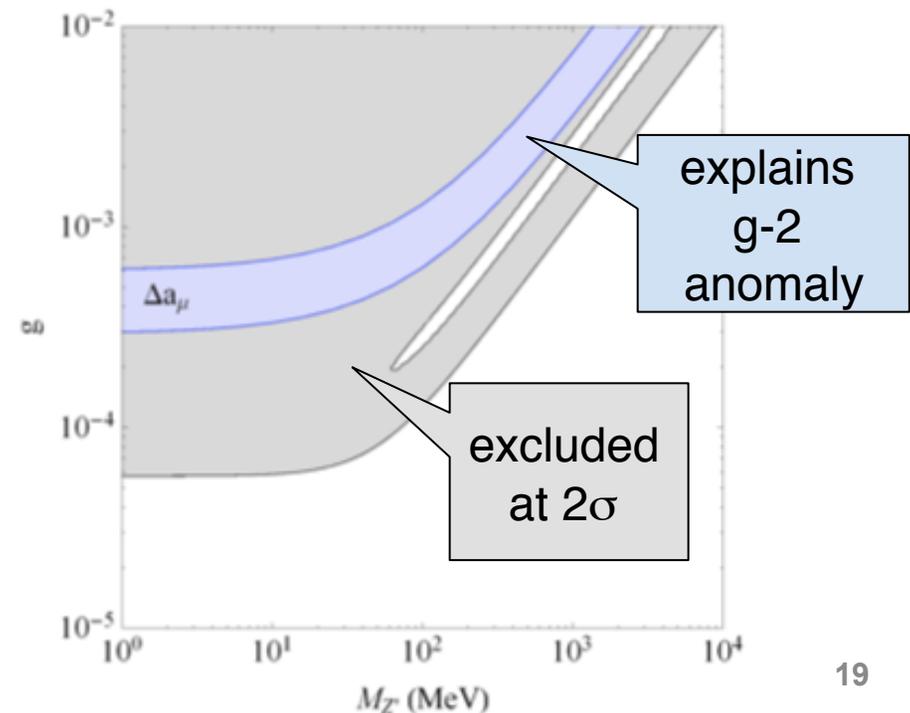
$$Q_{\alpha,SM}^2 = (Zg_p^V + Ng_n^V)^2 \quad \rightarrow \quad Q_{\alpha,NSI}^2 = \left[Z \left(g_p^V + \frac{3g^2}{2\sqrt{2}G_F(Q^2 + M_{Z'}^2)} \right) + N \left(g_n^V + \frac{3g^2}{2\sqrt{2}G_F(Q^2 + M_{Z'}^2)} \right) \right]^2$$

- Q^2 -dependence \rightarrow affects recoil spectrum
- 2 parameters: $g, M_{Z'}$



Dashed: SM
Solid: NSI w/ $M_{Z'} = 10 \text{ MeV}, g=10^{-4}$

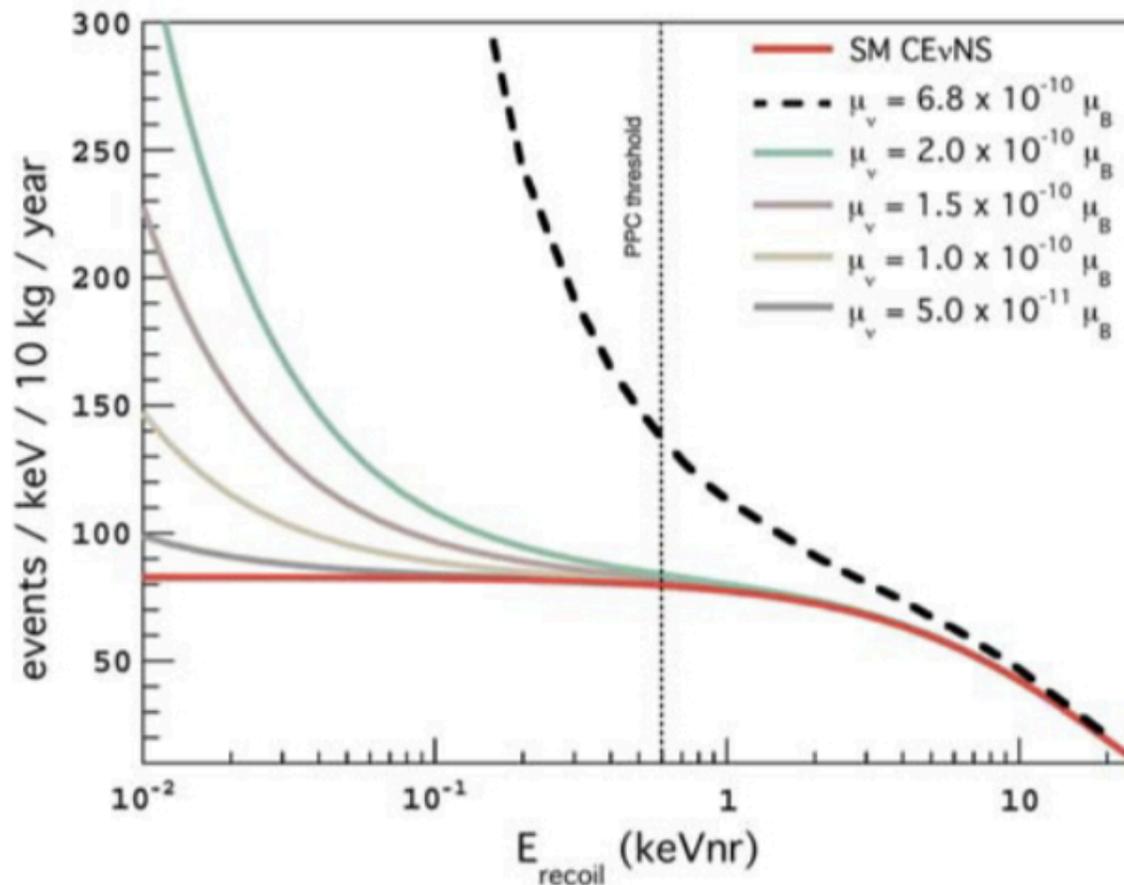
Blue: ν_μ
Red: $\nu_\mu + \bar{\nu}_\mu$
Black: $\nu_\mu + \bar{\nu}_\mu + \nu_e$



Neutrino magnetic moment

Signature is **distortion at low recoil energy E**

$$\left(\frac{d\sigma}{dT}\right)_m = \frac{\pi\alpha^2\mu_\nu^2 Z^2}{m_e^2} \left(\frac{1 - T/E_\nu}{T} + \frac{T}{4E_\nu^2}\right)$$



→ requires very low energy threshold (i.e., Ge)

See also Kosmas et al., arXiv:1505.03202

More in Juan's talk

Nuclear physics with CEvNS

If systematics can be reduced to ~ few % level,
we can start to explore nuclear form factors

P. S. Amanik and G. C. McLaughlin, J. Phys. G 36:015105

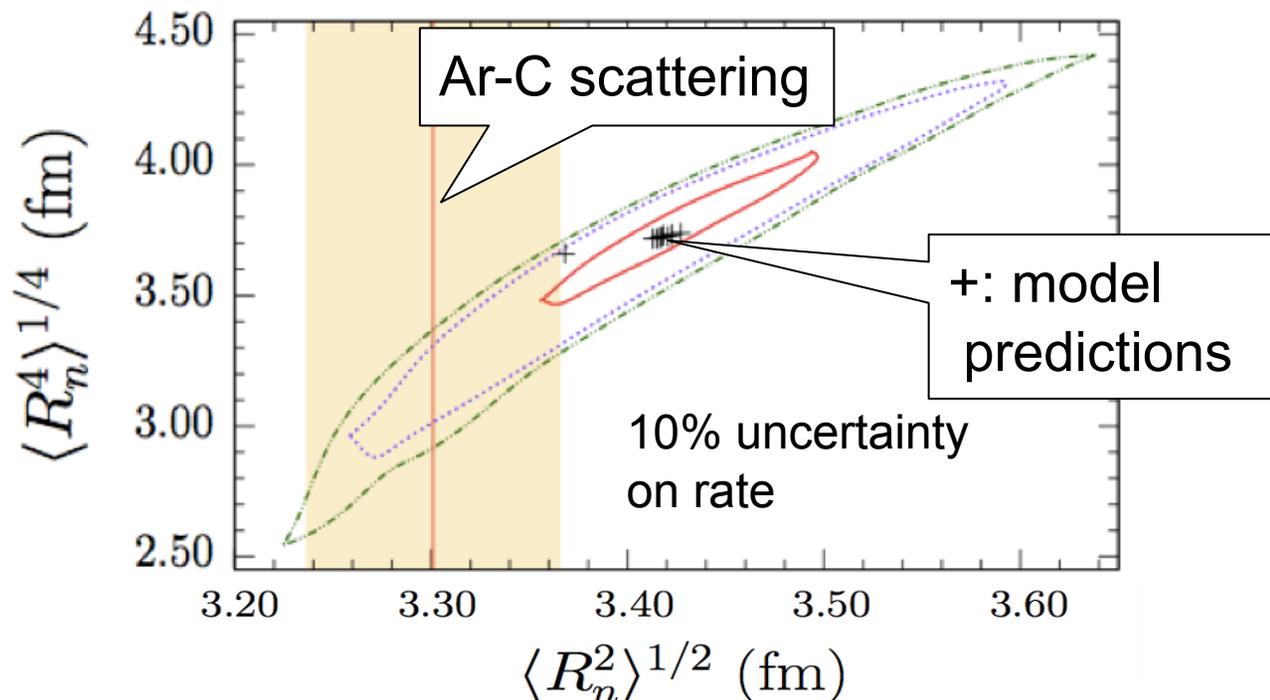
K. Patton et al., PRC86 (2012) 024612

$$\frac{d\sigma}{dT} = \frac{G_F^2 M Q_W^2}{2\pi \cdot 4} F^2(Q) \left(2 - \frac{MT}{E_\nu^2} \right)$$

Form factor: encodes information about nuclear (primarily neutron) distributions

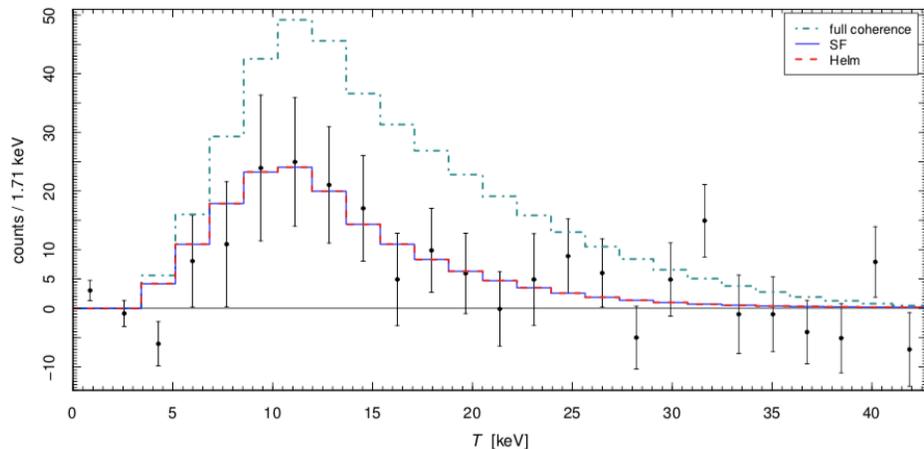
Fit recoil **spectral shape** to determine the $F^2(Q)$ moments
(requires very good energy resolution, good systematics control)

Example:
tonne-scale
experiment
at π DAR source



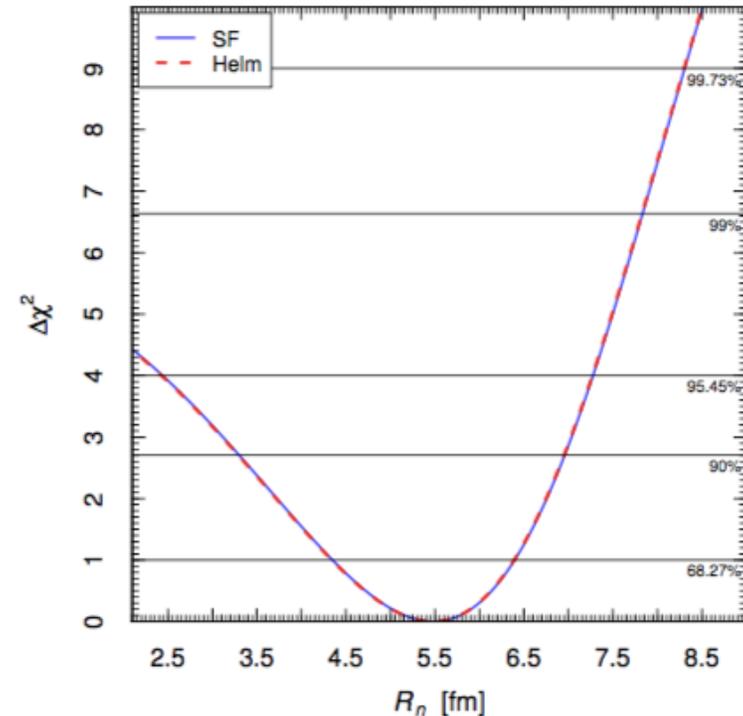
Sensitivity to R_n in the recoil spectrum shape

M. Cadeddu, C. Giunti, Y. F. Li, and Y. Y. Zhang. “Average CsI neutron density distribution from COHERENT data.” (2017). 1710.02730.



$$F_N^{\text{Helm}}(q^2) = 3 \frac{j_1(qR_0)}{qR_0} e^{-q^2 s^2/2},$$

$$R_n = 5.5^{+0.9}_{-1.1} \text{ fm.} \quad \Delta R_{np} \simeq 0.7^{+0.9}_{-1.1} \text{ fm.}$$



- Fit to neutron radius w/ $\sim 18\%$ uncertainty, but does not handle bin-by-bin correlation of systematics
- Also some info on neutron skin

More in Rex's talk

Light DM direct detection possibilities

Light new physics in coherent neutrino-nucleus scattering experiments

Patrick deNiverville,¹ Maxim Pospelov,^{1,2} and Adam Ritz¹

¹Department of Physics and Astronomy, University of Victoria, Victoria, BC V8P 5C2, Canada

²Perimeter Institute for Theoretical Physics, Waterloo, ON N2J 2W9, Canada

(Dated: May 2015)

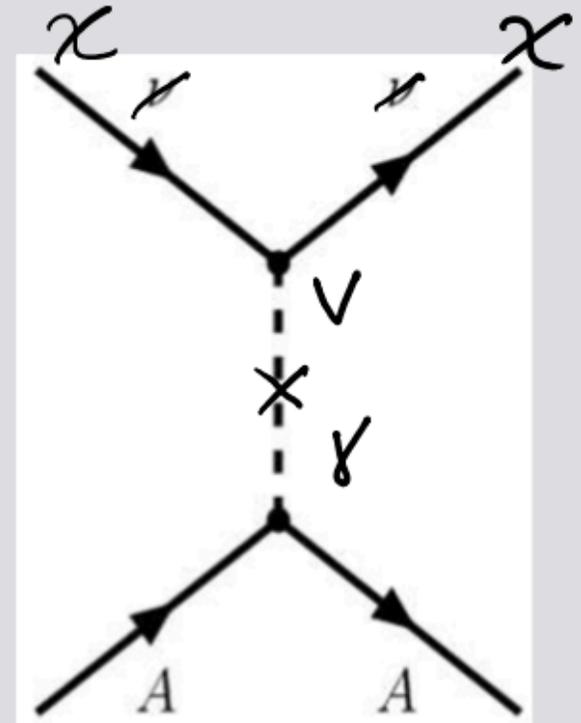
production:

proton \rightarrow target $\rightarrow \pi^{0,\pm} \rightarrow$

$$\pi^0 \rightarrow \gamma + V^{(*)} \rightarrow \gamma + \chi^\dagger + \chi$$

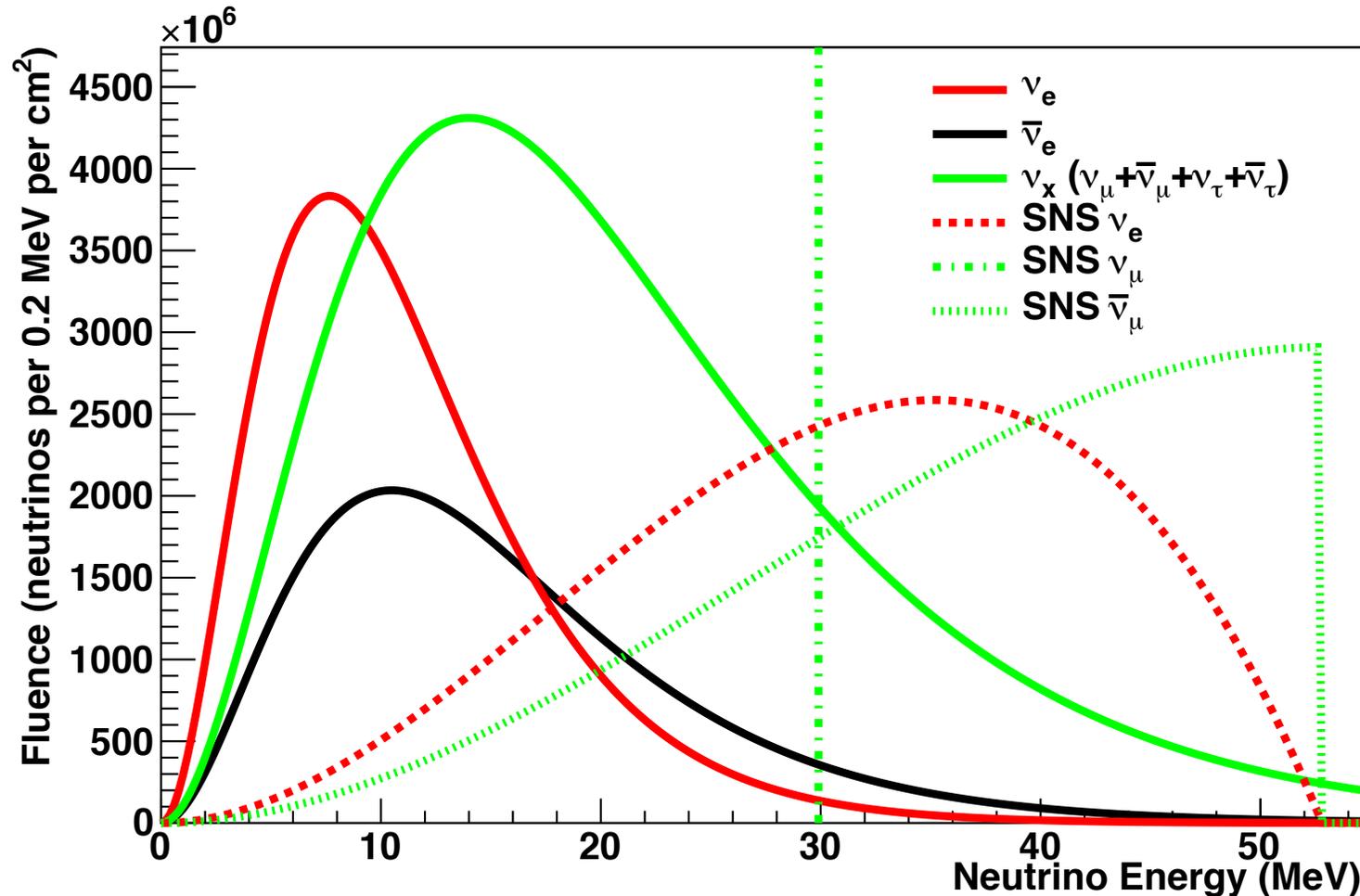
$$\pi^- + p \rightarrow n + V^{(*)} \rightarrow n + \chi^\dagger + \chi$$

detection:



More in Rex's talk

Non-CEvNS measurements for supernova neutrinos

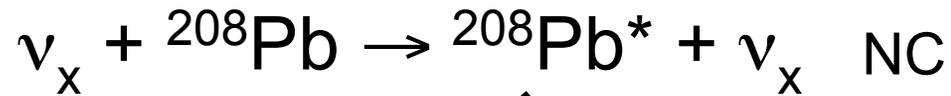


SNS spectrum overlaps well...
opportunities for measuring CC and NC

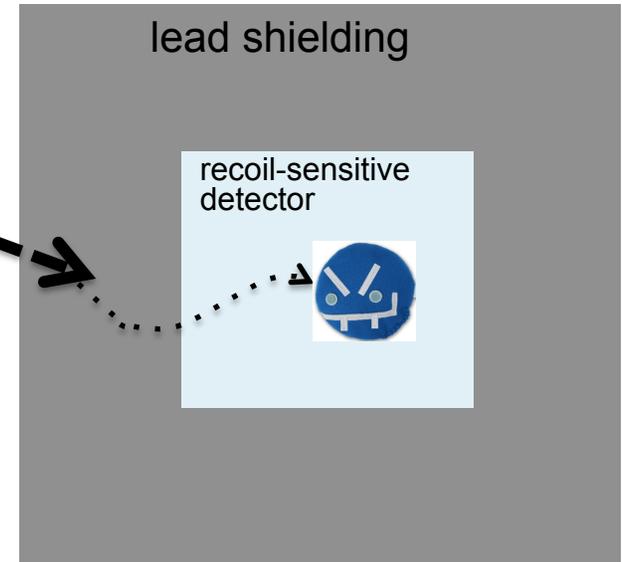
A “friendly fire” in-time background: Neutrino Induced Neutrons (NINs)



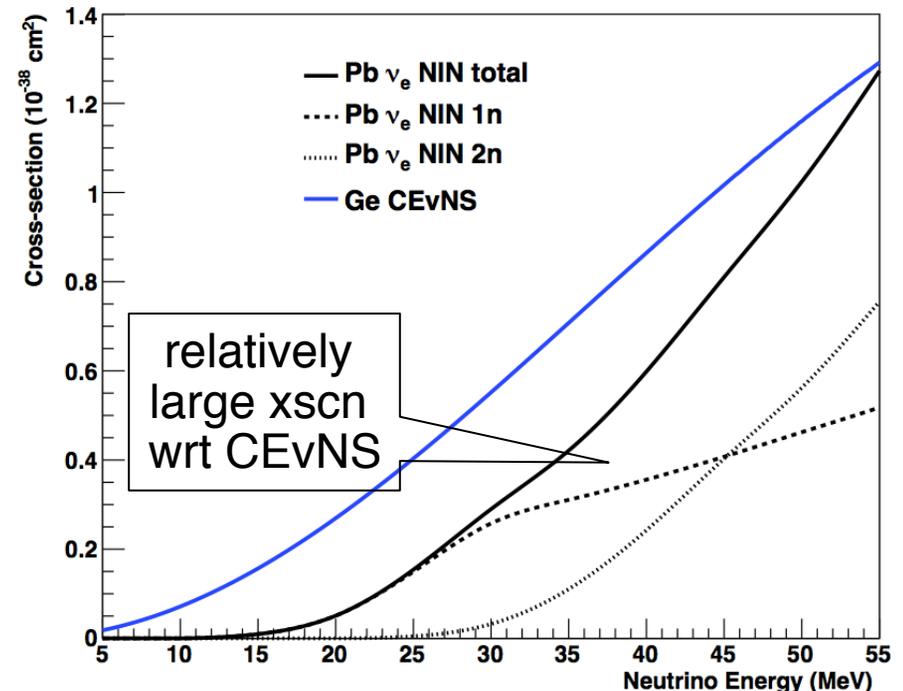
↓
1n, 2n emission



↓
1n, 2n, γ emission



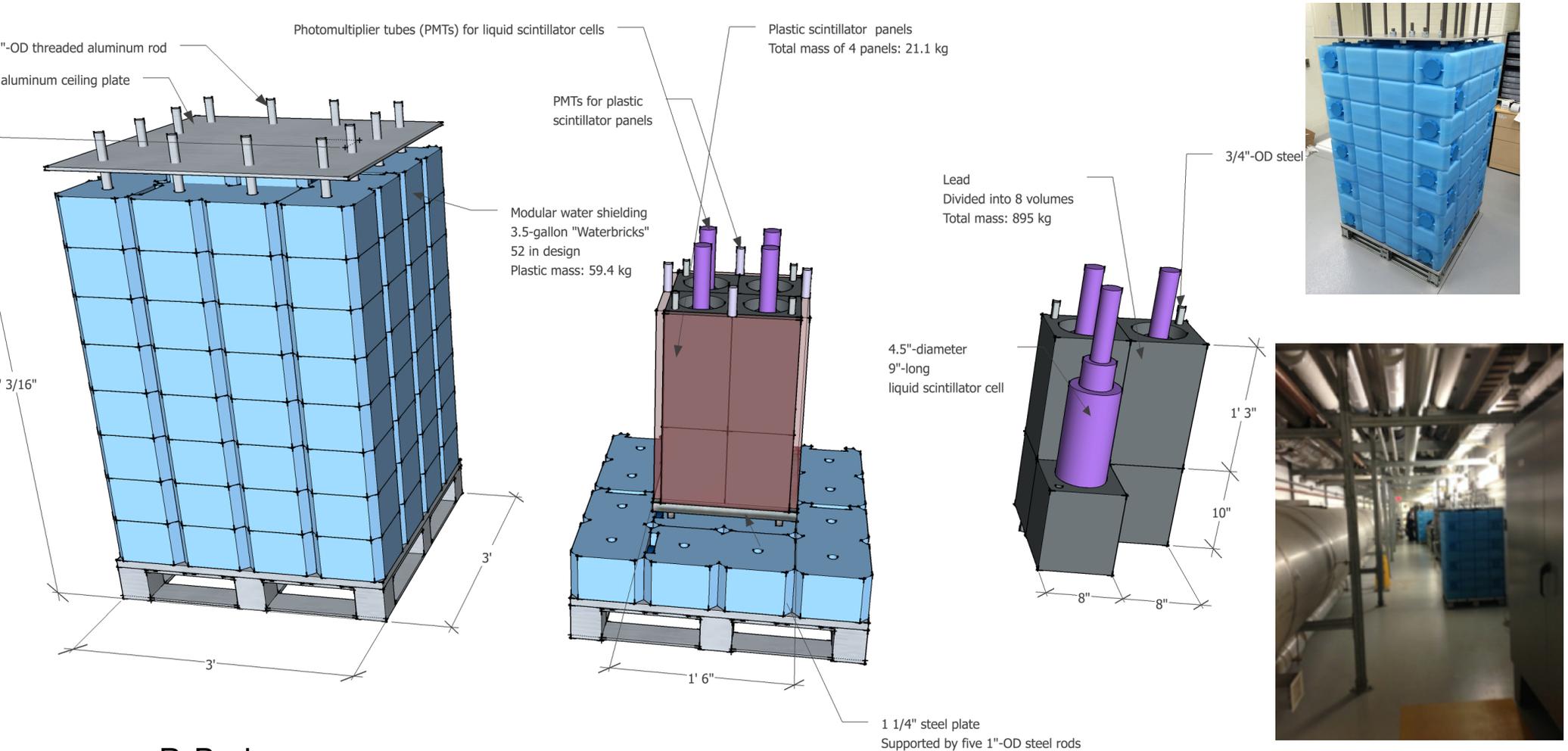
- potentially non-negligible background from shielding
- requires careful shielding design
- large uncertainties (factor of few) in xscn calculation
- [Also: a signal in itself, e.g, HALO SN detector]



NIN measurement in SNS basement with Nubes

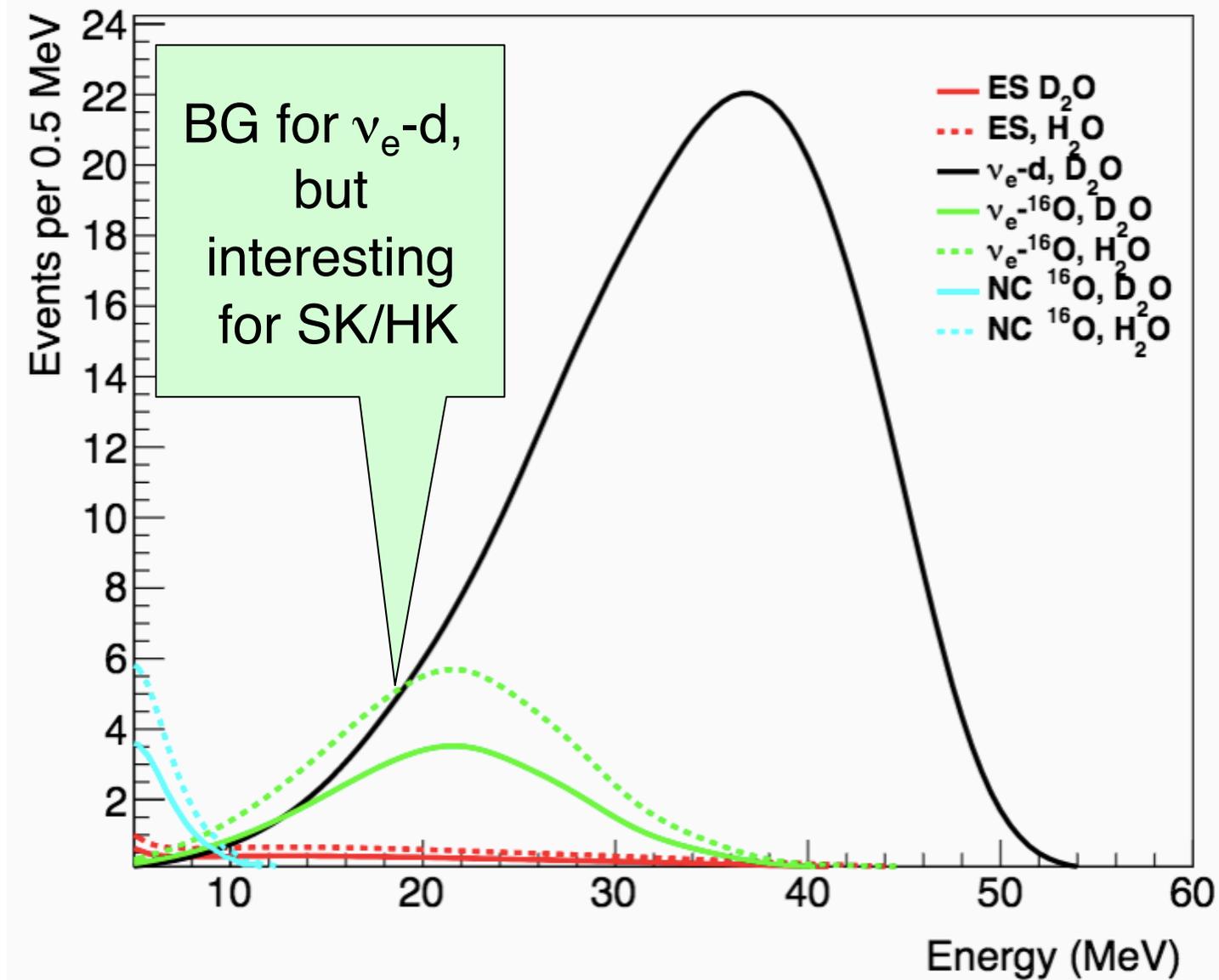
Liquid scintillator surrounded by Pb, Fe (swappable for other NIN targets)
inside water shield

Planning upgrade using PROSPECT Li-loaded scintillator



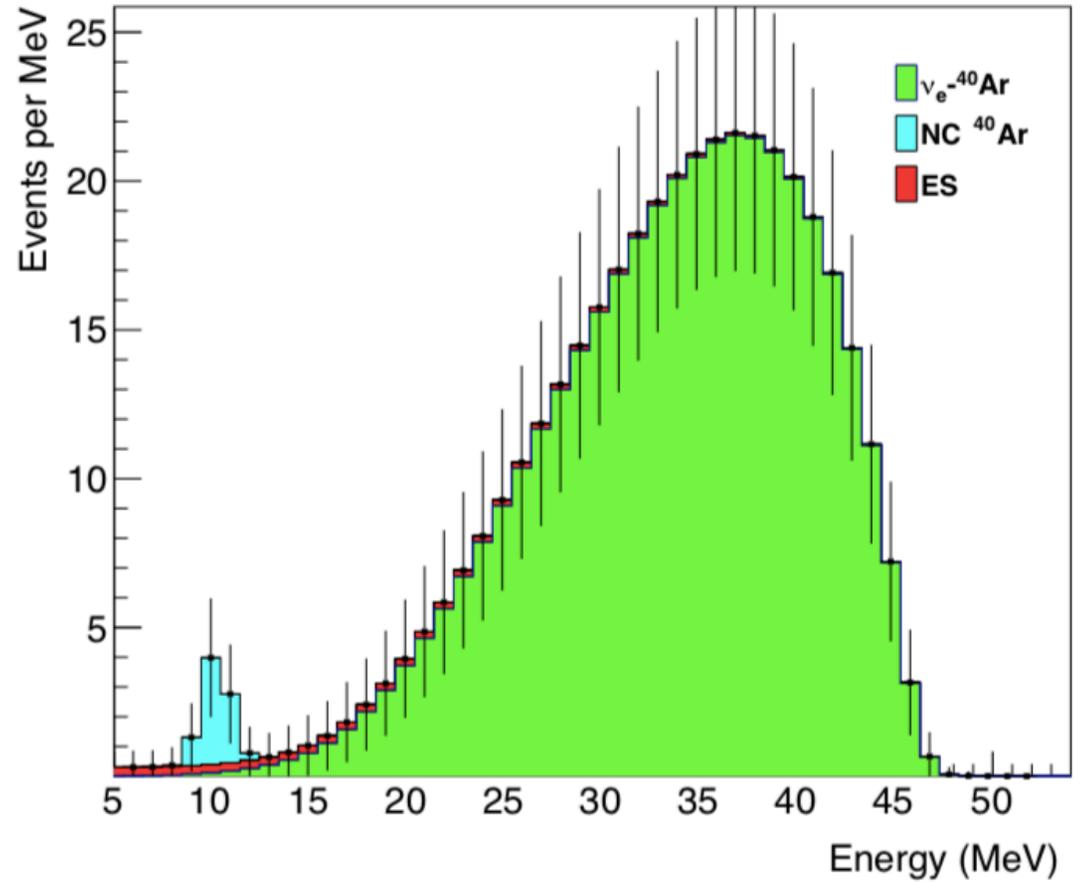
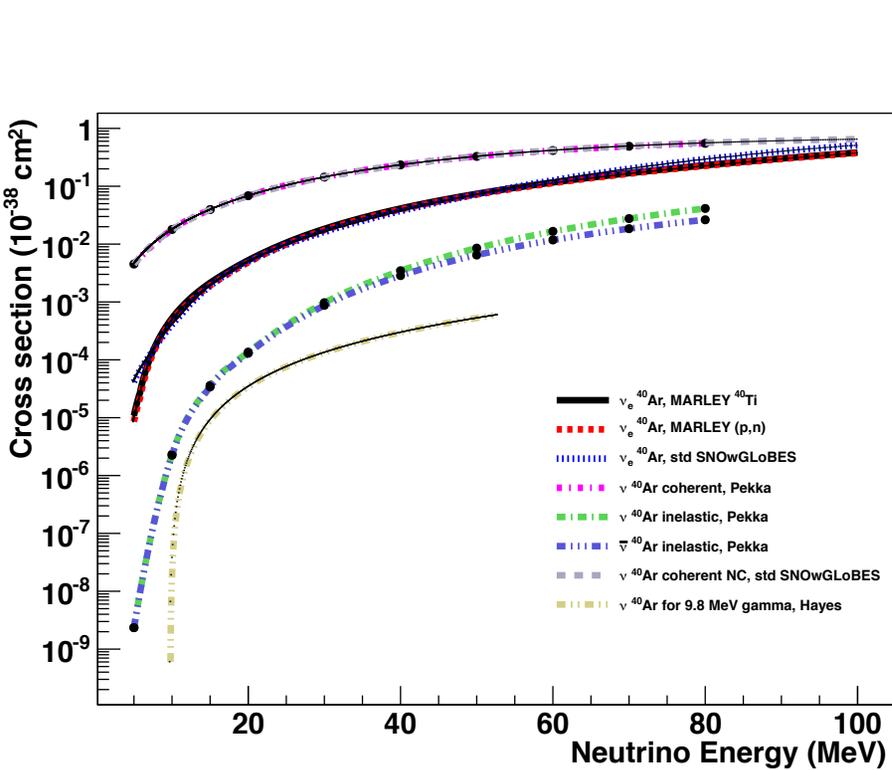
P. Barbeau

CC and NC measurements in light & heavy water



More in Jason's talk

CC & NC measurements in LAr



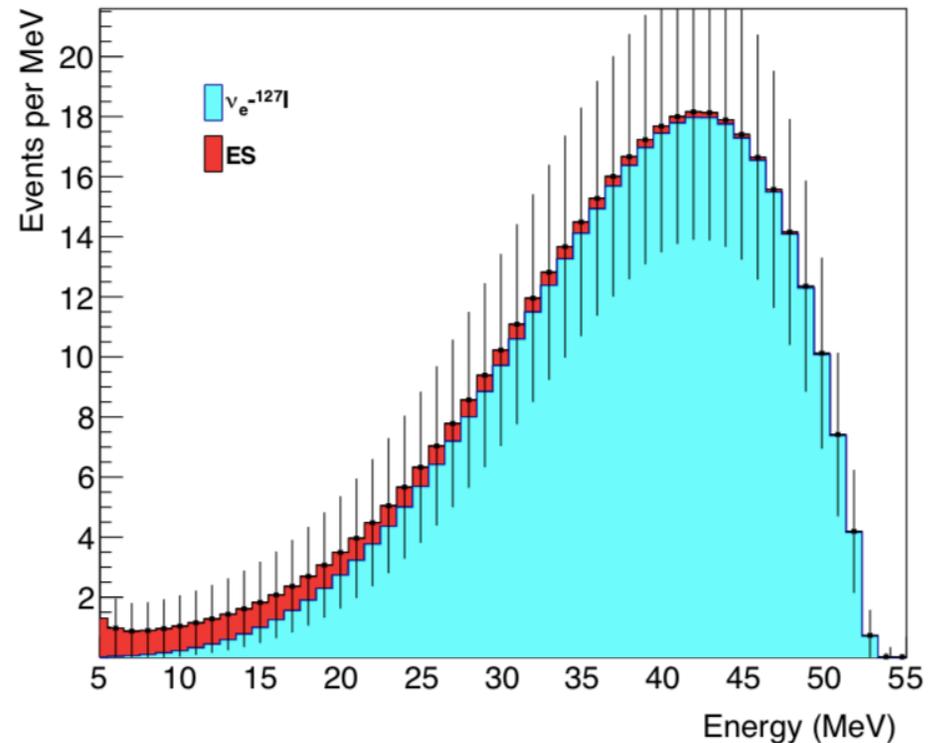
More in Rex's talk

Charged-current measurements in ^{127}I

TABLE III. Contributions of individual multipoles to the total cross section for neutrinos from muon decay, in units of 10^{-40} cm^2 . The two columns correspond to quenched and free values for g_A , respectively (see text).

J^π	$g_A = -1.0$	$g_A = -1.26$
0^+	0.096	0.096
0^-	0.00001	0.00002
1^+	1.017	1.528
1^-	0.006	0.008
2^+	0.155	0.213
2^-	0.693	1.055
3^+	0.149	0.171
3^-	0.017	0.025
Total	2.098	3.096

J. Engel, 1994



Exclusive cross section to bound final states of ^{127}Xe
 measured @ LANL, but we can measure **inclusive CC** xscn in NaI

More in Phil's talk

Reducing systematic uncertainties

2017 CsI measurement

Uncertainties on signal and background predictions	
Event selection	5%
Quenching factor	25%
Flux	10%
Form factor	5%
Total uncertainty on signal	28%
Beam-on neutron background	25%

Dominant uncertainty
(detector-dependent)

Next largest uncertainty
(affects all detectors)

- ancillary quenching factor measurements are important for the physics program
- D₂O for flux normalization also needed



P5 Science Drivers

- Use the Higgs boson as a new tool for discovery
- Pursue the physics associated with neutrino mass 
- Identify the new physics of dark matter 
- Understand cosmic acceleration: dark energy and inflation
- Explore the unknown: new particles, interactions, and physical principles. 

COHERENT connects to 3/5 of these

Recommendation 4: Maintain a program of projects of all scales, from the largest international projects to mid- and small-scale projects.

In addition,
it's a “small-scale”
project

Take-Away Messages

- COHERENT will demonstrate CEvNS N^2 **dependence**
- **Multiple targets** needed for this demonstration, as well as to test for BSM physics

- Significant **NSI sensitivity improvements** expected
 - important for interpretation of LBL oscillation experiments
 - *combination* of targets helps

- Many other physics topics : CC/NC for SNB, magnetic moment, etc.
- Control of systematics will require:
 - D₂O for flux
 - ancillary QF measurements